# Semilinear periodic equation with arbitrary nonlinear growth and data measure: mathematical analysis and numerical simulation 

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#### Abstract

In this work, we are interested in the existence, uniqueness, and numerical simulation of weak periodic solutions for some semilinear elliptic equations with data measures and with arbitrary growth of nonlinearities. Since the data are not very regular and the growths are arbitrary, a new approach is needed to analyze these types of equations. Finally, a suitable numerical discretization scheme is presented. Several numerical examples are given which show the robustness of our algorithm.


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## 1. Introduction

Periodic equations play an important role in the development of mathematical analysis of differential and partial differential equations. These problems appear in the modeling of many real-world phenomena, including fluid mechanics, pseudo-plastic flows, chemical reactions (resistivity of materials), nerve impulses (Fitzhugh-Nagumo problem), population dynamics (Lotka-Volterra system), combustion, morphogenesis, genetics, etc. Hundreds of articles on periodic problems have been published in various journals and conference proceedings, although there are still more questions than answers. We refer the reader to $[1-10]$ for a good introduction to periodic problems. These references contain review articles on ordinary periodic differential equations, which focus on the mathematical modeling of nonlinear equations and expose different solving methods. Among them are degree theory, variational methods, compactness methods, monotone methods, lower and upper solutions techniques, etc.

The purpose of this paper is to conduct a mathematical analysis and a numerical simulation of weak solutions for the semilinear equation with periodic boundary conditions.

Consider the following model equation

$$
\left\{\begin{array}{l}
u(t)-u^{\prime \prime}(t)+j(t, u(t))=f \text { in } \quad(0, T),  \tag{1}\\
u(0)=u(T), \quad u^{\prime}(0)=u^{\prime}(T),
\end{array}\right.
$$

where $T>0$ is a period, $j:[0, T] \times \mathbb{R} \rightarrow[0,+\infty[$ is measurable continuous function with respect to $u$, T-periodic with respect to $t$, i.e., $j(0, r)=j(T, r) \forall r$ (it allows to expand $j$ into a continuous periodic function on $\mathbb{R}$, by $j(t, r+k T)=j(t, r) \forall r \in(0, T))$ and $f$ is a given positive bounded Radon measure on $] 0, T$ [, T-periodic in the sense of the following definition.
Definition 1. We denote by $\mathcal{M}_{B}^{+}(0, T)$ the set of positive bounded Radon measures on $] 0, T$.
$f \in \mathcal{M}_{B}^{+}(0, T)$ is said to be T-periodic if there exists $f_{\varepsilon} \in C([0, T])^{+}$such that $f_{\varepsilon}(0)=f_{\varepsilon}(T)$ and

$$
\forall \phi \in C([0, T]),\langle f, \phi\rangle=\lim _{\varepsilon \rightarrow 0} \int_{0}^{T} f_{\varepsilon}(t) \phi(t) d t
$$

An example of a Radon measure to be 1-periodic is $f=\delta_{\frac{1}{2}}$ since the Lorentzian sequence:

$$
\begin{equation*}
f_{\varepsilon}(t)=\frac{1}{\pi \varepsilon} \frac{1}{1+\frac{(t-1 / 2)^{2}}{\varepsilon^{2}}} \tag{2}
\end{equation*}
$$

is 1-periodic continuous (in the sense that $f_{\varepsilon}$ is defined on $[0,1]$ by (1), and its extension beyond $] 0,1[$ is given by $f_{\varepsilon}(t+k)=f_{\varepsilon}(t)$ with $k \in \mathbb{Z}$ and $t \in(0,1)$. One can prove that $f_{\varepsilon}$ is convergent in the sense of measure to $\delta_{\frac{1}{2}}$.

The case: $j \equiv 0$ corresponds to the linear periodic problem. It has been widely studied in the literature due to the regularity of $f$. When $f$ is T-periodic and $f \in C([0, T])$ and $r \longrightarrow j(t, r)$ is globally Lipschitz, Coster et al. [10] proves the existence of a periodic solution $u \in W^{1,2}([0, T])$. Takemura et al. [11] considered the case where $f$ is 1-periodic and $f \in L^{2}(0, T)$ and they prove existence and uniqueness of a periodic solution $u \in H^{2}(0,1)$, in addition $u$ is expressed as

$$
u(t)=\int_{0}^{T} G(t-s) f(s) d s \quad(0<t<T)
$$

where $G$ is the Green function given by

$$
G(t-s)= \begin{cases}\frac{1}{T} \frac{e^{\frac{(t-s)}{T}}-1}{\left(e^{-1}-1\right)^{2}}\left(1+\left(e^{-1}-1\right) \frac{(t-s)}{T}\right) & (0<s \leqslant t<T), \\ \frac{1}{T} \frac{e^{\frac{(t-s)}{T}}}{\left(e^{-1}-1\right)^{2}}\left(e^{-1}+\left(e^{-1}-1\right) \frac{(t-s)}{T}\right) & (0<t<s<T) .\end{cases}
$$

In the case where $j$ actually depends on $t$ and $u$, i.e. $j=j(t, u(t))$, the problem is said to be semilinear. It has been analyzed by Ciarlet et al. [12], by using an optimization method and under the following assumptions: $f \in C([0, T]), r \rightarrow j(t, r)$ is differentiable nondecreasing and $\forall t \in[0, T],\left|\frac{\partial j(t, r)}{\partial r}\right|$ is bounded on the bounded set of $\mathbb{R}$.

In the present work, we are particularly interested in cases where $f$ is irregular and the growth of $j$ with respect to $u$ is arbitrary. Obviously, classical methods fail to prove the existence and new techniques must be used. We describe some of them here.

The other analysis that we deal with in this paper, is the simulation of the periodic solution of (1). Several methods for numerical analysis and simulation of periodic equations have been proposed in the literature. One of the numerical methods is the collocation method, see [13,14]. Samoilenko [15] proposed another quasi-linear numerical method. Here we will present the complete discretization of equation (1) by finite differences. Then we reduce the search for a periodic solution to the solution of a nonlinear system whose dimension is the number of nodes of the considered mesh. We then develop an algorithm based on the Newton-Raphson method to numerically simulate a large system and obtain an approximation of our periodic solution.

The rest of this paper is organized as follows. In Section 2, we present the exact problem statement and main results. In Section 3, we give the existence proof for the semilinear problem, if $f \in L^{2}(0, T)$. In Section 4, we construct an approximate problem for (1) with regular data whose existence will be a consequence of the previous section. After performing a priori estimations, we pass to the limit in the approximated problem and prove the main existence result. The last section is devoted to numerical simulation of our general problem. After proposing a numerical scheme based on finite differences, we present several numerical examples to demonstrate the efficiency and robustness of our proposed algorithm.

## 2. Statement of the main theoretical result

Throughout this paper we assume:
$\left.A_{1}\right) f \in \mathcal{M}_{B}^{+}(0, T)$ T-periodic (in the sens of Definition 1);
$\left.A_{2}\right) j:[0, T] \times \mathbb{R} \rightarrow[0,+\infty[$ a mesurable T-periodic function;
$\left.A_{3}\right) \forall t, r \rightarrow j(t, s)$ is continuous and nondecreasing and $j(t, 0)=0$;
$\left.A_{4}\right) \forall r \in \mathbb{R}, j(t, r) \in L^{1}(0, T)$.

Consider for $1 \leqslant p \leqslant \infty$,

$$
W_{p e r}^{1, p}(0, T)=\left\{u \in W^{1, p}(0, T), \text { such that } u(0)=u(T)\right\}
$$

equipped with the norm induced by $W^{1, p}(0, T)$

$$
\|u\|_{1, p}=\|u\|_{p}+\left\|u^{\prime}\right\|_{p} .
$$

In the case $p=2$, this space is noted by $H_{p e r}^{1}(0, T)$.
Now we introduce the notion of weak periodic solution of the problem (1) used here.
Definition 2. A function $u$ is said to be a weak T-periodic solution of the problem (1), if

$$
\left\{\begin{array}{l}
u \in W_{p r e}^{1,1}(0, T)  \tag{3}\\
\int_{0}^{T} u(t) \phi(t) d t+\int_{0}^{T} u^{\prime}(t) \phi^{\prime}(t) d t+\int_{0}^{T} j(t, u(t)) \phi(t) d t=\langle f, \phi\rangle \text { for all } \phi \in W_{p e r}^{1, \infty}(0, T) .
\end{array}\right.
$$

## Remark 1.

i) for all $1 \leqslant p \leqslant \infty, W_{p e r}^{1, p}(0, T) \subset C([0, T])$ with compact injection.
ii) $\langle$,$\rangle denotes the duality bracket between \mathcal{M}_{B}(0, T)$ and $L^{\infty}(0, T)$.
iii) if $u \in W_{\text {per }}^{1, \infty}(0, T)$, since $j$ satisfy $\left(A_{4}\right)$, then $j(t, u(t)) \in L^{1}(0, T)$, therefore all terms in (3) make sense.
Till the end of this paper, we denote by $C$ every generic and positive constant. We have the following main result.
Theorem 1. Assume that $\left(A_{2}\right)-\left(A_{4}\right)$ holds. Then for all $f \in \mathcal{M}_{B}^{+}(0, T) T$-periodic, there exists a weak nonnegative $T$-periodic solution $u$ of (1).

## 3. An auxiliary existence result

Consider $f \in L^{2}(0, T)$. One can obtain the following result.
Theorem 2. Let $f \in L^{2}(0, T)$ be T-periodic and $j$ satisfy $\left(A_{2}\right)-\left(A_{4}\right)$. Then there exists a unique nonnegative weak $T$-periodic solution of the problem

$$
\left\{\begin{array}{l}
u \in H_{p e r}^{1}(0, T)  \tag{4}\\
\int_{0}^{T} u(t) \phi(t) d t+\int_{0}^{T} u^{\prime}(t) \phi^{\prime}(t) d t+\int_{0}^{T} j(t, u(t)) \phi(t) d t=\int_{0}^{T} f(t) \phi(t) d t \text { for all } \phi \in H_{p e r}^{1}(0, T) .
\end{array}\right.
$$

In addition, if $f \geqslant 0$ then $u(t) \geqslant 0 \forall t \in[0, T]$.
Proof. Let us define the functional

$$
J: \begin{aligned}
& H_{p e r}^{1}(0, T) \rightarrow \mathbb{R} \\
& v \rightarrow \frac{1}{2} \int_{0}^{T}|v(t)|^{2} d t+\frac{1}{2} \int_{0}^{T}\left|v^{\prime}(t)\right|^{2} d t+\int_{0}^{T} J_{p}(t, v(t)) d t-\int_{0}^{T} f(t) v(t) d t,
\end{aligned}
$$

where $J_{p}(t, r)=\int_{0}^{r} j(t, s) d s$. Since $J_{p}$ and $\|u\|_{1,2}^{2}$ are convex then $J$ is convex. Now we will prove that $J$ is lower semi-continous. Consider for $C \in \mathbb{R}$, the set

$$
A=[J \leqslant C]=\left\{v \in H_{p e r}^{1}(0, T) \text { such that } J(v) \leqslant C\right\} .
$$

We are going to prove that $A$ is a closed set in $H_{p e r}^{1}(0, T)$. Let us consider a sequence $v_{n} \in A$ and $v_{n} \rightarrow v$ in $H_{p e r}^{1}(0, T)$, we have

$$
\begin{equation*}
\frac{1}{2} \int_{0}^{T}\left|v_{n}(t)\right|^{2} d t+\frac{1}{2} \int_{0}^{T}\left|v_{n}^{\prime}(t)\right|^{2} d t+\int_{0}^{T} J_{p}\left(t, v_{n}(t)\right) d t-\int_{0}^{T} f(t) v_{n}(t) d t \leqslant C . \tag{5}
\end{equation*}
$$

Since $H_{p e r}^{1}(0, T) \subset C([0, T])$ with a compact injection, we can extract a subsequence $v_{n k}$ such that

$$
v_{n k} \rightarrow v \text { in } C([0,1]),
$$

since $v \in H_{p e r}^{1}(0, T)$ we get also $\int_{0}^{T} f v_{n_{k}} \rightarrow \int_{0}^{T} f v$ and by using Fatou's Lemma, we get

$$
\int_{0}^{T} J_{p}(v(t)) d t \leqslant \liminf _{k \rightarrow+\infty} \int_{0}^{T} J_{p}\left(t, v_{n}(t)_{k}\right) d t
$$

Passing to the limit in (5), we obtain $J(v) \leqslant \liminf _{n \rightarrow+\infty} J\left(v_{n k}\right) \leqslant C$. Therefore $v \in A$.

Now we prove that $J$ is infinite at infinity. We have

$$
J(v) \geqslant \frac{1}{2}\|u\|_{1,2}^{2}-\|v\|_{1,2}\|f\|_{L^{2}}
$$

then

$$
\liminf _{n \rightarrow+\infty} \frac{J(v)}{\|u\|_{1,2}}=+\infty
$$

Consequently, $J$ attains a unique global minimum

$$
\inf _{v \in H_{p e r}^{1}(0, T)} J(v)=\min _{v \in H_{p e r}^{1}(0, T)} J(v)=J(u)
$$

Let us finally show that $u$ is a solution of (4). By choosing $v=u+s \phi$ for any $s$ in the neighborhood of 0 and any $\phi \in H_{p e r}^{1}(0, T)$, we get:

$$
\begin{aligned}
\frac{1}{2} \int_{0}^{T}|u(t)+s \phi(t)|^{2} d t & +\frac{1}{2} \int_{0}^{T}\left|u^{\prime}(t)+s \phi^{\prime}(t)\right|^{2} d t+\int_{0}^{T} J_{p}(t, u(t)+s \phi(t)) d t \\
& -\frac{1}{2} \int_{0}^{T}|u(t)|^{2} d t-\frac{1}{2} \int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t-\int_{0}^{1} J_{p}(t, u(t)) d t \geqslant s \int_{0}^{T} f(t) \phi(t) d t
\end{aligned}
$$

We divide the inequality by $s>0$, then $s<0$, the limit when $s$ approaches 0 gives us:

$$
\lim _{s \rightarrow 0} \frac{J(u+s \phi)-J(u)}{s}=0
$$

then

$$
\left.\frac{d}{d s} \right\rvert\, s=0 J(u+s \phi)=0
$$

Which, in turn, yields

$$
\int_{0}^{T} u(t) \phi(t) d t+\int_{0}^{T} u^{\prime}(t) \phi^{\prime}(t) d t+\int_{0}^{T} j(t, u(t)) \phi(t) d t=\int_{0}^{T} f(t) \phi(t) d t \quad \forall \phi \in H_{p e r}^{1}(0, T)
$$

Finally, suppose $f \geqslant 0$ a.e. in $(0, T)$, since $j$ is nonnegative, we consider the equation (1) with

$$
\hat{j}(t, r)=\left\{\begin{array}{l}
j(t, r) \text { if } r \geqslant 0 \\
0 \text { if } r<0
\end{array}\right.
$$

instead of $j$. It is clear that if $r \geqslant 0, \hat{j}=j$.
We introduce the function sign ${ }^{-}$defined on $\mathbb{R}$ by

$$
\operatorname{sign}^{-}(r)=\left\{\begin{array}{r}
-1 \text { if } r<0 \\
0 \text { if } r \geqslant 0
\end{array}\right.
$$

as sign ${ }^{-}$is an increasing function, we consider the convex function $\rho_{\varepsilon}$, which is a twice differentiable function such that

$$
\rho_{\varepsilon}^{\prime}(r) \rightarrow \operatorname{sign}^{-} r \text { when } \varepsilon \rightarrow 0
$$

Let us take $\rho_{\varepsilon}^{\prime}(u)$ as a test function, then, we get

$$
\int_{0}^{T} u(t) \rho_{\varepsilon}^{\prime}(u(t)) d t+\int_{0}^{T} u^{\prime 2}(t) \rho_{\varepsilon}^{\prime \prime}(u(t)) d t+\int_{0}^{T} \hat{j}(t, u(t)) \rho_{\varepsilon}^{\prime}(u(t))=\int_{0}^{T} \rho_{\varepsilon}^{\prime}(u(t)) f(t) d t
$$

using the convexity of $\rho_{\varepsilon}$, we deduce that

$$
\int_{0}^{T} u^{\prime 2}(t) \rho_{\varepsilon}^{\prime \prime}(u(t)) d t \geqslant 0
$$

for the other terms, we have

$$
\lim _{\varepsilon \rightarrow 0} \int_{0}^{T} \hat{j}(t, u(t)) \rho_{\varepsilon}^{\prime}(u(t)) d t=\lim _{\varepsilon \rightarrow 0} \int_{[u \geqslant 0]} \hat{j}(t, u(t)) \rho_{\varepsilon}^{\prime}(u(t)) d t+\int_{[u<0]} \hat{j}(t, u(t)) \rho_{\varepsilon}^{\prime}(u(t)) d t
$$

It follows that

$$
=\int_{[u<0]} \hat{j}(t, u) d t=0
$$

$$
\lim _{\varepsilon \rightarrow 0} \int_{0}^{T} u(t) \rho_{\varepsilon}^{\prime}(u(t)) d t \leqslant \lim _{\varepsilon \rightarrow 0} \int_{0}^{T} \rho_{\varepsilon}^{\prime}(u(t)) f(t) d t
$$

which implies that

$$
\int_{0}^{T} u^{-}(t) d t \leqslant-\int_{0}^{T} f(t) d t \leqslant 0
$$

which allows us to conclude that $u(t) \geqslant 0$ a.e. $t$ in $[0, T]$.

## 4. Proof of the main result

Since $f \in \mathcal{M}_{B}^{+}(0, T)$ T-periodic then there exist $f_{n} \in C([0, T]) f_{n}(0)=f_{n}(1), f_{n} \geqslant 0$ such that $\left\|f_{n}\right\|_{L^{1}} \leqslant\|f\|_{\mathcal{M}_{B}}$ and which converge to $f$ in $\mathcal{M}_{B}^{+}(0, T)$. According to the Theorem 2, there exists $u_{n}$, nonnegative weak T-periodic solution of

$$
\left\{\begin{array}{l}
u_{n} \in H_{p e r}^{1}(0, T), u_{n} \geqslant 0  \tag{6}\\
\int_{0}^{T} u_{n}(t) \phi(t) d t+\int_{0}^{T} u_{n}^{\prime}(t) \phi^{\prime}(t) d t+\int_{0}^{T} j_{n}\left(t, u_{n}(t)\right) \phi(t) d t=\int_{0}^{T} f_{n}(t) \phi(t) d t \quad \forall \phi \in H_{p e r}^{1}(0, T)
\end{array}\right.
$$

We have the following estimates.
Lemma 1. Let $u_{n}$ be the sequence defined as above, then we have:
i) $\int_{0}^{T}\left|u_{n}\right| d t \leqslant\|f\|_{\mathcal{M}_{B}}$;
ii) $\int_{0}^{T}\left|j_{n}\left(t, u_{n}\right)\right| d t \leqslant\|f\|_{\mathcal{M}_{B}}$;
iii) $\int_{0}^{T}\left|u_{n}^{\prime \prime}(t)\right| \leqslant C\|f\|_{\mathcal{M}_{B}}$.

Proof. Take $\phi \equiv 1$ in (6), and as $j\left(\cdot, u_{n}\right) \geqslant 0$, it comes that

$$
\int_{0}^{T} u_{n}(t) d t+\int_{0}^{T} j_{n}\left(t, u_{n}(t)\right) d t=\int_{0}^{1} f_{n}(t) d t \leqslant\|f\|_{\mathcal{M}_{B}}
$$

since $u_{n}$ and $j\left(t, u_{n}\right) \geqslant 0$, then we obtain i) and ii).
Finally, we deduce from (6) that

$$
\left\{\begin{array}{l}
u_{n}^{\prime \prime}=u_{n}+j\left(t, u_{n}\right)-f_{n} \text { in } \mathcal{D}^{\prime}(0, T)  \tag{7}\\
u_{n}(0)=u_{n}(T) \\
u_{n}^{\prime}(0)=u_{n}^{\prime}(T)
\end{array}\right.
$$

Since $u_{n}, j\left(t, u_{n}\right), f_{n} \in L^{1}(0,1)$, then $u_{n}^{\prime \prime} \in L^{1}(0, T)$ and one get

$$
\int_{0}^{T}\left|u_{n}^{\prime \prime}(t)\right| d t \leqslant C\|f\|_{\mathcal{M}_{B}}
$$

which proves iii).
Furthermore, $u_{n}^{\prime}$ is continue and $u_{n}^{\prime}(0)=u_{n}^{\prime}(T)$, then there exists $t_{0 n}$ such that $u_{n}^{\prime}\left(t_{0 n}\right)=0$, hence $u_{n}^{\prime}(t)=\int_{t_{0 n}}^{t} u_{n}^{\prime \prime}(s) d s$. According to ii) of Lemma 1, we get $\int_{0}^{T}\left|u_{n}^{\prime}(t)\right| d t \leqslant C\|f\|_{\mathcal{M}_{B}}$. Then $u_{n}$ is bounded in $W_{\text {per }}^{1,1}(0, T)$.

Since $W_{p e r}^{1,1}(0, T) \subset C([0, T])$ with compact injection, then there exists $u \in W_{p e r}^{1,1}(0, T)$ and a subsequence noted by $u_{n}$ such that $u_{n} \rightarrow u$ in $C[0,1]$. Therefore, due to $\left(A_{3}\right), j\left(t, u_{n}\right) \rightarrow j(t, u)$ in $L^{1}(0, T)$.

This allows us to go to the limit in the equation (6) and obtain that $u$ is a weak periodic solution of the equation (1).

## 5. Numerical simulation

In this section, we propose a numerical simulation of the equation (1) using finite differences. The first subsection is devoted to discretizing our periodic problem using the finite difference method, and then we present a solution algorithm based on the Newton-Raphson method. In the second subsection, we show numerical results obtained depending on the case where the source $f$ is a regular function or a Radon measure.

### 5.1. Discretization and numerical algorithm

For that we discrete the interval $[0, T]$ in $N+1$ points $x_{k}=(k-1) * h$, for $k=1, \ldots, N+1$, where $h=\frac{T}{N}$. Let us set $u_{k}=u\left(x_{k}\right)$ and add two fictious points $x_{0}=-h, x_{N+2}=T+h$. Let us denote $u_{0}=u\left(x_{0}\right)$ and $u_{N+2}=u\left(x_{N+2}\right)$. Therefore, since $u$ is periodic we have $u_{N+1}=u_{1}, u_{0}=u_{N}$ and $u_{N+2}=u_{2}$, we then have $N$ unknowns $u_{k}, k=1,2, \ldots, N$.

Our problem (1) can then be discretized in space as follows:

$$
\left\{\begin{array}{l}
u_{1}-\frac{1}{h^{2}}\left(u_{N}-2 u_{1}+u_{2}\right)+j\left(0, u_{1}\right)=f(0), \\
u_{i}-\frac{1}{h^{2}}\left(u_{i+1}-2 u_{i}+u_{i-1}\right)+j\left(x_{i}, u_{i}\right)=f\left(x_{i}\right) \text { for } 2 \leqslant i \leqslant N-1, \\
u_{N}-\frac{1}{h^{2}}\left(u_{1}+u_{N-1}-2 u_{N}\right)+j\left(x_{N}, u_{N}\right)=f\left(x_{N}\right) .
\end{array}\right.
$$

This can be written in matrix form:

$$
\begin{equation*}
G(U)=\left(\mathcal{I}_{N}-\frac{1}{h^{2}} A\right) * U+J(U)-F=0 \tag{8}
\end{equation*}
$$

where $U=\left(u_{i}\right)_{1 \leqslant i \leqslant N}$ is the unknows vector, $F=\left(f\left(x_{i}\right)\right)_{1 \leqslant i \leqslant N}, \mathcal{I}_{N}$ is the identity matrix of order $N$, the matrix $A$ is given by

$$
\mathbf{A}=\left[\begin{array}{cccccc}
-2 & 1 & 0 & \cdots & 0 & 1 \\
1 & -2 & 1 & 0 & \cdots & 0 \\
& \ddots & \ddots & \ddots & & \\
& & \ddots & \ddots & \ddots & \\
0 & \cdots & 0 & 1 & -2 & 1 \\
1 & 0 & \cdots & 0 & 1 & -2
\end{array}\right]
$$

and the nonlinear term vector $J(U)$ is given by

$$
J(U)=\left[\begin{array}{c}
j\left(0, u_{1}\right) \\
\cdots \\
j\left(x_{i}, u_{i}\right) \\
\ldots \\
j\left(x_{N}, u_{N}\right)
\end{array}\right]
$$

We will use the Newton-Raphson method to solve equation (8) starting from the initial $U_{0}$ which is the solution of the linear system:

$$
\begin{equation*}
\left(\mathcal{I}_{N}-\frac{1}{h^{2}} A\right) * U_{0}=F \tag{9}
\end{equation*}
$$

Our algorithm is therefore the following.

```
Algorithm 1
    Input: choose \(k_{\max }\) the maximum number of iterations, the tolerance \(\varepsilon_{0}\), we get \(N+1\) points \(x_{i}=(i-1) * h\),
    \(h=\frac{T}{N}\)
    set \(k=0\) and set \(U=U_{0}\)
    repeat
        set \(k=k+1\)
        if ( \(k=k_{\text {max }}\) ) then exit convergence
        endif
        set \(Y=G(U)\)
        solve \(\left(\mathcal{I}_{N}-\frac{1}{h^{2}} A+D J(U)\right) * D=-Y\)
        set \(U=U+D\)
    until \(\|D\|<\varepsilon_{0}\)
    Output: \(U\)
```


### 5.2. Numerical examples

We present some numerical examples depending on the cases if the source $f$ is regular or not. The first example is the following

$$
\left\{\begin{array}{l}
u(t)-u^{\prime \prime}(t)+u(t)^{4}=1+t \sin (\pi t) \text { in }(0,1),  \tag{10}\\
u(0)=u(1), \quad u^{\prime}(0)=u^{\prime}(1) .
\end{array}\right.
$$

The simulation we give here corresponds to $T=1, \varepsilon_{0}=1 e^{-9}, N=400, k_{\max }=8$.
Figure 1 shows the shape of the periodic solution and Figure 2 shows the decrease of the norm between two successive iterations as a function of the iteration number.


Fig. 1. Shape of the periodic solution of (10).


Fig. 2. Error of Newton.

The second example is the following

$$
\left\{\begin{array}{l}
u(t)-u^{\prime \prime}(t)+u^{4}(t)=\delta_{\frac{1}{2}} \text { in }(0,1),  \tag{11}\\
u(0)=u(1), \quad u^{\prime}(0)=u^{\prime}(1) .
\end{array}\right.
$$

We have approximated the Dirac mass $\delta_{\frac{1}{2}}$ by the sequence of Lorentzian 1-periodic function $\left(f_{\varepsilon}\right)$ we defined before (2). The simulation we give here corresponds to $T=1, \varepsilon=1 e^{-12}, \varepsilon_{0}=1 e^{-9}, N=500$, $k_{\text {max }}=8$.

Figure 3 shows the shape of the periodic solution and Figure 4 shows the decrease of the norm between two successive iteration as a function of the iteration number.


Fig. 3. Shape of the periodic solution of (11).


Fig. 4. Error of Newton.

## 6. Conclusions

In this work, we have been interested in the mathematical analysis and numerical simulation of a class of periodic nonlinear equations with non-regular data. If the data is regular, we prove the existence and uniqueness of the periodic solution through optimization methods. With the data only nonnegative measure, we construct a sequence of periodic solutions based on the regular case, and after obtaining a priori estimates, we show that we can extract a subsequence that converges to the solution to the problem we consider. We then propose a numerical algorithm to simulate these periodic solutions, giving some examples of when the data is regular or irregular. Numerical simulations demonstrate that our algorithm is efficient and robust. Besides, in the future we will focus on analyzing other numerical methods for simulating periodic equations, such as FEM (Finite Element Method), ANN (Artificial Neural Networks), LBM (Lattice Boltzmann Method) etc. and analyzed the performance differences between these methods, comparing their accuracy, time consumption, etc.

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# Напівлінійне періодичне рівняння з довільною нелінійністю зростання та мірою даних: математичний аналіз та чисельне моделювання 

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У цій роботі цікавимося існуванням, єдиністю та чисельним моделюванням слабких періодичних розв'язків для деяких напівлінійних еліптичних рівнянь із мірами даних та з довільними нелінійностями зростання. Оскільки дані не дуже регулярні, а зростання є довільним, необхідний новий підхід для аналізу цих типів рівнянь. Накінець, наведено відповідну чисельну схему дискретизації. Наведено декілька числових прикладів, які демонструють надійність запропонованого алгоритму.

Ключові слова: періодичний розв'язок; напівлінійне рівняння; метод оптимізаціӥ; чисельне моделювання.

